

Investigation of intersubband transitions in wide bandgap oxide quantum well structures for optoelectronic device applications A. Atić^{1,2*}, N. Vuković², J. Radovanović²

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MOTIVATION

- □ There is a growing interest in wide bandgap oxides like ZnO, MgO, Ga₂O₃, etc. [1], [2].
- □ These materials possess unique properties suitable for highperformance optoelectronic devices in sensing, communications, and imaging applications.
- \Box Ga₂O₃ has a wide bandgap energy of about 4.8 to 4.9 eV, making it transparent to UV radiation and opaque to visible light [3].
- □ ZnO exhibits exceptional optical and electrical properties, including a high exciton binding energy and substantial oscillator strength [4].
- □ Intersubband transitions in multiple quantum well (MQW) structures hold promise for efficient light absorption and emission in the mid-infrared to terahertz spectral range.

SIMULATION RESULTS

 \Box We study a multiple quantum well design, focusing on Mg_xZn_{1-x}O/ZnO and $(Al_xGa_{1-x})_2O/Ga_2O_3$ materials. The results of our numerical simulations are absorption profiles for these material systems.



□ The structure shown is for the Mg_{0.3}Zn_{0.7}O/ZnO material system. The same layer thickness is further used for simulating the

□ The study aims to numerically simulate the absorption spectra of wide bandgap oxide MQW structures, focusing on the light-matter interaction where the dominant many-body contribution is the depolarization field.

THEORETICAL MODEL

□ 1D effective mass Schrödinger equation:

$$-\frac{\hbar^2}{2}\frac{d}{dz}\frac{1}{m^*(z)}\frac{d\psi_i(z)}{dz} + U_{eff}(z)\psi_i(z) = E_i\psi_i(z),$$

$$U_{eff}(z) = U_c(z) - e\varphi(z) + U_{xc}(z) - eF_e z$$

 \Box U_{xc} is the local exchange correlation potential:

$$U_{xc}(z) = -\frac{e^4}{32\pi^2\hbar^2} \frac{m^*(z)}{\varepsilon^*(z)^2} \left(\frac{9\pi}{4}\right)^{\frac{1}{3}} \frac{2}{\pi r_s^*} \left[1 + 0.0545r_s^* \log\left(1 + \frac{11.4}{r_s^*}\right)\right]^{\frac{1}{3}}$$

□ For the electrostatic potential, the Poisson equation reads:



□ The single particle absorption coefficient is [5]:





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$$\alpha_{2D,s}(\omega) = C_s \sum_{\alpha} f_{\alpha} \Delta N_{\alpha} G(\omega - \omega_{\alpha})$$

The Hamiltonian describing the intersubband plasmon:

$$T_{plasmon} = \sum_{\alpha} \hbar \widetilde{\omega}_{\alpha} p_{\alpha}^{\dagger} p_{\alpha} + \frac{\hbar}{2} \sum_{\alpha \neq \beta} \Xi_{\alpha,\beta} (p_{\alpha} + p_{\alpha}^{\dagger}) (p_{\beta} + p_{\beta}^{\dagger})$$

□ The coupling due to dipole-dipole Coulomb interaction is described by the coupling strength: $\Xi_{\alpha,\beta} = \frac{\omega_{P\alpha}\omega_{P\beta}}{2\sqrt{\widetilde{\omega}_{P\alpha}\widetilde{\omega}_{P\beta}}} \frac{S_{\alpha\beta}}{\sqrt{S_{\alpha\alpha}S_{\beta\beta}}}, \qquad S_{\alpha\beta} = \frac{1}{\hbar\omega_{\alpha}} \frac{1}{\hbar\omega_{\beta}} \left(\frac{\hbar^2}{2m^*}\right)^2 \int_{-\infty}^{+\infty} dz \xi_{\alpha}(z)\xi_{\beta}(z),$

$$\xi_{\alpha}(z) \equiv \xi_{ij}(z) = \psi_i(z) \frac{\partial \psi_j(z)}{\partial z} - \psi_j(z) \frac{\partial \psi_i(z)}{\partial z}.$$

 \Box The new N frequency W_N can be calculated by diagonalizing the following $2N \times 2N$ matrix: 0 1 $\left[\widetilde{\omega}\right]_{\alpha}$

$$M = \begin{bmatrix} I_1 & C_{12} & \dots & C_{1N} \\ C_{12} & I_2 & \dots & C_{2N} \\ \vdots & \vdots & \ddots & \vdots \\ C_{1N} & C_{2N} & \dots & I_N \end{bmatrix}, \qquad I_{\alpha} = \begin{bmatrix} \varpi_{\alpha} & -\widetilde{\omega}_{\alpha} \\ 0 & -\widetilde{\omega}_{\alpha} \end{bmatrix},$$

Current density for n-th multisubband plasmon:

$$e\hbar \qquad \nabla \xi_{\alpha}(z) \sqrt{\Delta N_{\alpha}}$$

□ The layer sequence in nm, starting from the left barrier reads: 3/3.1/2.5/2.4/3.4/5.5/3. The widest well is uniformly dopped with $n = 30 \times 10^{18} cm^{-3}$

 \Box The electronic structure is calculated at applied field of F= 7.3 \times 10⁶ $\frac{V}{2}$

REFERENCES

- [1] J. Y. Tsao et al., Adv. Electron. Mater., vol. 4, no. 1, p. 1600501, Jan. (2018). [2] A. V. Osipov, S. S. Sharofidinov, E. V. Osipova, A. V. Kandakov, A. Y. Ivanov, and S. A. Kukushkin, Coatings, vol. 12, no. 12, p. 1802, Nov. (2022).
- [3] J. E. Lyman and S. Krishnamoorthy, *J. Appl. Phys.*, vol. 127, no. 17, p. 173102, May (2020).
- [4] T. Makino et al., Appl. Phys. Lett., vol. 76, no. 24, pp. 3549–3551, Jun. (2000).
- [5] G. Pegolotti, A. Vasanelli, Y. Todorov, and C. Sirtori, Phys. Rev. B, vol. 90, no. 3, p. 035305, Jul (2014).
- [6] V. P. Sirkeli and H. L. Hartnagel, Opto-Electronics Rev., vol. 27, no. 2, pp. 119–122, Jun. (2019)

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• Absorption can be calculated by integrating all the current densities associated with different multisubband plasmons:



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